## Simple and Accurate Determination of Global $\tau_{\rm R}$ in Proteins Using <sup>13</sup>C or <sup>15</sup>N Relaxation Data

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Received September 1, 1999; revised December 10, 1999

In the study of protein dynamics by <sup>13</sup>C or <sup>15</sup>N relaxation measurements different models from the Lipari–Szabo formalism are used in order to determine the motion parameters. The global rotational correlation time  $\tau_R$  of the molecule must be estimated prior to the analysis. In this Communication, the authors propose a new approach in determining an accurate value for  $\tau_R$  in order to realize the best fit of  $R_2$  for the whole sequence of the protein, regardless of the different type of motions atoms may experience. The method first determines the highly structured regions of the sequence. For each corresponding site, the Lipari–Szabo parameters are calculated for  $R_1$  and NOE, using an arbitrary value for  $\tau_R$ . The  $\chi^2$  for  $R_2$ , summed over the selected sites, shows a clear minimum, as a function of  $\tau_R$ . This minimum is used to better estimate a proper value for  $\tau_R$ . @ 2000 Academic Press

*Key Words:* rotational correlation time; relaxation parameters; protein dynamics; residues selection; factorial discriminant analysis.

Recent advances in isotope labeling methods have enabled <sup>13</sup>C and <sup>15</sup>N nuclear spin relaxation to become widely used for studying internal motions in macromolecules (proteins). The relaxation data ( $R_1$ ,  $R_2$ , NOE) are usually obtained at a single value of the static magnetic field, allowing a poor sampling of the spectral density function. The simplest way to modelize this function is by means of the Lipari–Szabo formalism (1, 2), that is, depending on only three parameters:

$$J(\omega) = \frac{2}{5} \left[ \frac{S^2 \tau_{\rm R}}{1 + (\omega \cdot \tau_{\rm R})^2} + \frac{(1 - S^2) \tau_{\rm e}}{1 + (\omega \cdot \tau_{\rm e})^2} \right], \qquad [1]$$

where  $\tau_{\rm R}$  is the correlation time of the global motion (supposed isotropic) of the molecule in solution,  $\tau_{\rm e}$  is the effective correlation time of the internal motion, and  $S^2$  represents the spatial restriction of the internal motion.

The determination of the rotational correlation time constitutes a primary requirement in the attempt at motion analysis in all molecular sites of the protein, starting from the experimental relaxation data obtained from labeled nuclei.

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The standard procedure for estimating this parameter is to use the  $R_2/R_1$  ratio for selected sites in the protein sequence (3). This approach has been extensively discussed recently (4-6) and it is based on two principal ways of selecting residues.

One way requires that the selected sites should fulfill the extreme narrowing limit, which is a very restrictive condition. It assumes the second term of Eq. [1] to be negligible (7) (model 1 of LS formalism). The value for  $\tau_{\rm R}$  is then estimated from Eq. [2]:

$$\tau_{\rm R} = \frac{1}{2\omega_{\rm X}} \sqrt{\frac{6R_2}{R_1} - 7}.$$
 [2]

Due to the constraint imposed on the minimum value of  $\tau_e$ , which is seldom fulfilled,  $\tau_R$  is usually underestimated using this method (4). Besides, the proper sites might not be easily selectable a priori. Attempts have been made to improve the selection (8), based on NOE data, but nevertheless, the condition of independence of  $R_2/R_1$  on  $S^2$  still remains.

The second, less restrictive way to select residues consists of fitting simultaneously all three relaxation parameters with  $\tau_{\rm R}$ ,  $S^2$ , and  $\tau_{\rm e}$ . The determination of  $\tau_{\rm R}$  in the latter approach has been carried out using the local site treatment of Shurr *et al.* (9) or using the globally linked approach of Dellwo and Wand (10). In this case the three LS parameters ( $\tau_{\rm R}$ ,  $S^2$ , and  $\tau_{\rm e}$ ) are calculated by minimizing the  $\chi^2$  function described by

$$\chi^{2} = \sum_{\text{residues}} \left[ \frac{(R_{1\text{exp}} - R_{1\text{calc}})^{2}}{\sigma^{2}(R_{1})} + \frac{(\text{NOE}_{\text{exp}} - \text{NOE}_{\text{calc}})^{2}}{\sigma^{2}(\text{NOE})} + \frac{(R_{2\text{exp}} - R_{2\text{calc}})^{2}}{\sigma^{2}(R_{2})} \right]$$
[3]

in order to fit best, simultaneously, all three experimental relaxation parameters for each residue.

We propose a simpler method, in which the site selection is done using the general model 4 of the Lipari–Szabo formalism (7), accordingly fitting the relaxation data. In this way, for a given value for  $\tau_{\rm R}$ , one can obtain the values for  $S^2$  and  $\tau_{\rm e}$  in





**FIG. 1.**  $C^{\alpha}$  spin-spin relaxation rates as a function of residue index in the NCS protein.

Eq. [1] which fit exactly  $R_1$  and NOE (11). Then, the value for  $\tau_R$  is that obtained from the best fit of  $R_2$  among the selected sites, using only the exact estimates of  $S^2$  and  $\tau_e$  from  $R_1$  and NOE relaxation data. The extreme narrowing limit condition is no longer required but the selected sites should be submitted to fast and restrictive motion. They are usually located in the well-structured regions of the proteins and  $R_2$  is a sensible parameter for making this choice.

In the case of very restrictive motions (low amplitude,  $S^2 \approx$ 1), the corresponding time scale of the correlation time is usually in the range of picoseconds. For these nuclei, the relaxation parameters  $R_1$ ,  $R_2$ , and NOE can be fitted using Eq. [1] for  $J(\omega)$  in the Lipari–Szabo formalism. For other nuclei submitted to a slow motion regime in the time scale of nanoseconds, it is necessary to apply an extension of the Lipari-Szabo formalism (12). In these cases the value of  $R_2$  strongly diminishes. Moreover, there are cases where nuclei may have very slow motions contributing only to  $R_2$  relaxation mechanisms and not to  $R_1$  or NOE. For these specific residues, characterized by conformational changes on a time scale ranging in the domain of micro- to milliseconds, the  $R_2$  values are increasing. Such an example is shown in Fig. 1 for the NCS protein (13). In the figure, one may easily distinguish two regions in the sequence, characterized by a practically constant  $R_2$  value corresponding to highly structured domains in the protein. On the contrary, the edges of the sequence are characterized by decreased  $R_2$  values, indicating the presence of nuclei with fast and high amplitude motions, in the time scale of nanoseconds. For these specific cases, where motions are modeled on two widely separated time scales, the Lipari-Szabo formalism proves to be inappropriate and the extended model-free approach of Clore et al. is invoked (12). Finally one may notice on the plot a residue (83) where the  $R_2$  value is

strongly increasing, indicating conformational changes affecting only the spin-spin relaxation rate.

The Lipari–Szabo formalism is concerned only with very fast and hindered motions. Hence, it is very appropriate to fit the dynamic parameters with  $R_1$  and NOE that are insensible to slow motions. Assuming a known value for  $\tau_R$ , it is possible to find unique solutions for  $S^2$  and  $\tau_e$ , exactly fitting the experimental data for  $R_1$  and NOE (11).

We can suggest a new approach to determine  $\tau_R$  from the  $R_2$  fitting, in these highly structured domains of the sequence, where  $R_2$  tends to be practically constant. This method is based only on LS formalism, taking into account the characteristics of the internal motion of the nuclei, without any further assumption on  $S^2$  or  $\tau_e$ .

From the  $R_1$  and NOE fittings at an arbitrary chosen  $\tau_R$  value, one could reconstruct the  $R_2$  values at each site. The comparison with experimental data is done by calculating the corresponding  $\chi^2$  function,

$$\chi^2 = \sum_{\text{residues}} \frac{(R_{2\text{exp}} - R_{2\text{calc}})^2}{\sigma^2(R_2)},$$
[4]

where the sum is done over the selected sites. Note that, here,  $R_{1\text{calc}}$  and NOE<sub>calc</sub> equal exactly  $R_{1\text{exp}}$  and NOE<sub>exp</sub>, respectively, so that Eq. [3] reduces to Eq. [4].

The sites involved in the summation of Eq. [4] can be straightforwardly selected by examination of  $R_2$  (Fig. 1), but a more objective method is provided by multivariate statistics analysis. For example, a factorial discriminant analysis (14) performed on  $R_1$ ,  $R_2$ , and NOE experimental data allow us to clearly distinguish three regions (Fig. 2), of motionally similar sites. Residues inside region 2 have all three relaxation parameters fitting the LS formalism while those inside regions 1 and 3 are submitted to a slow motion regime (extension of LS) and exchange broadening, respectively.

The dependence of the  $\chi^2$  function over  $\tau_R$  presents a clear minimum (Fig. 3), which gives the best value for the rotational correlation time. The error on the estimation of  $\tau_R$  may be



**FIG. 2.** Factorial discriminant analysis on experimental relaxation data  $(R_1, R_2, \text{NOE})$  used for determining motional similarities among residues.



**FIG. 3.**  $\chi^2(R_2)$  calculated from the set  $(S^2, \tau_e)$  values, fitting  $R_1$  and NOE data for the arbitrary chosen value of  $\tau_R$  (the LS formalism, Eq. [1]), as a function of  $\tau_R$ .

obtained from a Monte-Carlo simulation. This method has been tested and compared (Fig. 4) with the  $R_2/R_1$  method (15), using data obtained for the NCS protein, totally enriched in <sup>13</sup>C, measuring relaxation rates of C $\alpha$ , at 500 MHz. Three different selections of residues were used. The first one, giving the smaller  $\chi^2$ , was restricted to a few residues belonging to the most structured regions of the protein. A second selection was extended to practically all of the protein backbone, except for the few residues at the edges and residue 83 which clearly exhibits an exchange contribution to  $R_2$ . This selection corresponds to all residues contained inside region 2 of Fig. 2. It gives the largest  $\chi^2$ , due to the larger number of terms in the summation (Eq. [4]). Finally a third selection was tested. It includes the residues inside the gray region (Fig. 2) of the factorial discrimant analysis map. Figure 4 clearly shows that the  $\tau_{\rm R}$  values obtained by the  $R_2/R_1$  method are systematically lower than those obtained by the method proposed in this paper, as expected (for none of the selected residues is the extreme narrowing limit fulfilled). Furthermore, it must be emphasized that the  $\tau_{\rm R}$  values thus obtained do not depend on the residue selection that is different from the  $R_2/R_1$  case which gives  $\tau_{\rm R}$  values dependent on the selected sites. As a result, the calculated value for  $\tau_{\rm R}$  using the sites inside region 2 of the FDA map is 4.53 ns ( $\sigma = 0.07$ ) for the NCS protein at 35°C.

We have also applied this method in order to obtain  $\tau_{\rm R}$  for the NCS protein for four temperatures (35, 40, 45, and 50°C). The variation of determined correlation times  $\tau_{\rm R}$  in function of  $\eta/T$  (Fig. 5) shows the expected linear dependence predicted by the Stokes–Einstein equation:

$$\tau_{\rm R} = \frac{4\pi\eta_{\rm w}r_{\rm H}^3}{3kT}.$$
 [5]

Finally, a global exchange process or an anisotropic motion may be a drawback of this approach but it has the advantages



**FIG. 4.**  $\chi^2(R_2)$  as a function of  $\tau_R$  in a Monte-Carlo simulation. To the left,  $\tau_R$  is calculated from Eq. [2]  $(R_2/R_1 \text{ method})$  using three different residue selections. To the right is the result obtained with the proposed method, using the same residue selections.



**FIG. 5.**  $\tau_{\rm R}$  dependence on the  $(\eta/T)$  term in the Stokes–Einstein equation as determined experimentally in NCS.

of rapid determination of motion parameters from measurements at a single value of the static magnetic field. It can also determine the proper  $\tau_{R}$  value using the minimum hypothesis, fitting the best possible motion parameters inside the LS formalism—model 4 (7)—for the highly structured residues of the protein.

## ACKNOWLEDGMENT

N.I.P. was supported by a postdoctoral fellowship from the Ligue Nationale Contre le Cancer (France).

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